

CLIQUE POTENTIALS 8

- Whatever factorization we ^pick, we know that only connected nodes can be arguments of ^a single local function.
- \bullet A *clique* \mathbf{x}_c is a fully connected subset of nodes.
- Thus, consider using ^a product of positive clique potentials:

$$
P(\mathbf{x}) = \frac{1}{Z} \prod_{\text{cliques } c} \psi_c(\mathbf{x}_c) \qquad Z = \sum_{\mathbf{x}} \prod_{\text{cliques } c} \psi_c(\mathbf{x}_c)
$$

- The product of functions that don't need to agree with each other.
- Still factors in the way that the graph semantics demand.
- Without loss of generality we can restrict ourselves to maximal cliques. (Why?)

• We often represent the clique potentials using their logs:

 $\psi_C(\mathbf{x}_C) = \exp\{-H_C(\mathbf{x}_C)\}\$

for arbitrary real valued "energy" functions $H_C(\mathbf{x}_C).$ The negative sign is ^a standard convention.

• This gives the joint ^a nice additive structure:

$$
P(\mathbf{x}) = \frac{1}{Z} \exp\{-\sum_{\text{cliques } C} H_C(\mathbf{x}_c)\} = \frac{1}{Z} \exp\{-H(\mathbf{x})\}
$$

where the sum in the exponent is called the "free energy":

$$
H(\mathbf{x}) = \sum_C H_C(\mathbf{x}_c)
$$

• This way of defining ^a probability distribution based on energies is the "Boltzmann distribution"from statistical physics.

Partition Function 11

- \bullet Normalizer $Z(\mathbf{x})$ above is called the "partition function".
- Computing the normalizer and its derivatives can often be the hardest part of inferene and learning in undirected models.
- Often the factored structure of the distribution makes it possible to efficiently do the sums/integrals required to compute Z .
- \bullet Don't *always* have to compute Z , e.g. for conditional probabilities.

 \bullet The model implies $\textbf{x} \perp \textbf{z} \mid \textbf{y}$

$$
p(\mathbf{x}, \mathbf{y}, \mathbf{z}) = p(\mathbf{y})p(\mathbf{x}|\mathbf{y})p(\mathbf{z}|\mathbf{y})
$$

• We can write this as:

$$
p(\mathbf{x}, \mathbf{y}, \mathbf{z}) = p(\mathbf{x}, \mathbf{y})p(\mathbf{z}|\mathbf{y}) = \psi_{\mathbf{xy}}(\mathbf{x}, \mathbf{y})\psi_{\mathbf{yz}}(\mathbf{y}, \mathbf{z})
$$

$$
p(\mathbf{x}, \mathbf{y}, \mathbf{z}) = p(\mathbf{x}|\mathbf{y})p(\mathbf{z}, \mathbf{y}) = \psi_{\mathbf{xy}}(\mathbf{x}, \mathbf{y})\psi_{\mathbf{yz}}(\mathbf{y}, \mathbf{z})
$$

cannot have all potentials be marginals cannot have all potentials be conditionals

• The positive clique potentials can only be thought of as general "compatibility", "goodness" or "happiness" functions over their variables, but not as probability distributions.

(a) (b) No directed model can represent these and only these independencies. $\mathbf{x} \perp \mathbf{y} \mid \{\mathbf{w}, \mathbf{z}\}$ $\mathbf{w} \perp \mathbf{z} \mid \{\mathbf{x},\mathbf{y}\}$

Z

• No.

No undirected model can represent these and only these independencies. x ⊥ y

Z

Hammersley-Clifford Theorem (1971) ¹³

- H-C theorem tells us that the family of distributions defined by the conditional independence semantics on the graph and the family defined by products of potential functions[∗] on maximal cliques are the same. $(*$ arbitrary real valued, but strictly positive)
- For directed models, there is ^a version of this theorem which tells us that the family of distributions defined by the conditional independencies semantics of the directed graph and the family defined by products of parent-conditionals are the same.
- Notice the crucial difference between *graphs*, which tells us independencies that are true no matter what local functions we choose, and numerical functions which could introduce some extra independencies, once we know them.

Example: Ising Models 15

- \bullet Common model for binary nodes: spin-glass/ Ising lattice.
- Nodes are arranged in ^a regular topology (often ^a regular packing grid) and connected only to their geometric neighbours.

- For example, if we think of each node as ^a ^pixel, we might want to encourage nearby pixels to have similar intensities.
- Energy is of the form:

$$
H(\mathbf{x}) = \sum_{ij} \beta_{ij} \mathbf{x}_i \mathbf{x}_j + \sum_i \alpha_i \mathbf{x}_i
$$

Example: Gaussian Distribution 16

- The most common and important undirected graphical model on ^a set of continuous valued nodes is the Gaussian (normal).
- It uses pairwise potentials between every pair of nodes to define an energy identical to the Ising model, but for continuous values:

$$
H(\mathbf{x}) = \sum_{ij} (\mathbf{x}_i - \mu_i) V_{ij} (\mathbf{x}_j - \mu_j)
$$

where μ is the mean and V is the inverse covariace matrix.

• Like ^a "fully connected" lattice.

Also, the Gaussian is the maximum entropy distribution consistent with the mean and covariance defined by μ and ${\bf V}$.

Example: Boltzmann Machines 17

• Fully observed Boltzmann machines are the binary equivalent of ^a Gaussian distribution: fully connected Ising models on ^a set of binary random variables. (Also maxent.) Energy is the same:

 $H(\mathbf{x}) = \sum_{i,j} \beta_{ij} \mathbf{x}_i \mathbf{x}_j + \sum_{i} \alpha_i \mathbf{x}_i$

- •
- Boltzmann machines also add the possibility of having some units (random variables) which are never observed. These are called "hidden units" or "latent variables" and we will see much more about them later.
- For continuous variables, the equivalent of ^a Boltzmann machine with hidden units is called ^a "factor analysis" model.